

ENEE 420
FALL 2010
COMMUNICATIONS SYSTEMS
ANGLE MODULATION

Throughout, we consider the information-bearing signal $m : \mathbb{R} \rightarrow \mathbb{R}$. Its Fourier transform is given by

$$M(f) := \int_{\mathbb{R}} m(t) e^{-j2\pi ft} dt, \quad f \in \mathbb{R}.$$

Frequency modulation _____

The FM waveform $s_{\text{FM}} : \mathbb{R} \rightarrow \mathbb{R}$ associated with the information-bearing signal m is given by

$$s_{\text{FM}}(t) = A_c \cos(\theta_{\text{FM}}(t)), \quad t \in \mathbb{R}$$

with

$$\theta_{\text{FM}}(t) = 2\pi f_c t + 2\pi k_F \int_0^t m(r) dr, \quad t \in \mathbb{R}.$$

Phase modulation _____

The PM waveform $s_{\text{PM}} : \mathbb{R} \rightarrow \mathbb{R}$ associated with the information-bearing signal m is given by

$$s_{\text{PM}}(t) = A_c \cos(\theta_{\text{PM}}(t)), \quad t \in \mathbb{R}$$

with

$$\theta_{\text{PM}}(t) = 2\pi f_c t + k_P m(t), \quad t \in \mathbb{R}.$$

Single-tone modulating signals _____

In an attempt to understand how the spectrum of angle-modulated signals is shaped by that of the modulating signal, we consider the simple case of a single-tone modulating signal $m : \mathbb{R} \rightarrow \mathbb{R}$, say

$$m(t) = A_m \cos(2\pi f_m t), \quad t \in \mathbb{R}$$

with amplitude $A_m > 0$ and frequency $f_m > 0$. In that case, we note that

$$\begin{aligned} \theta_{\text{FM}}(t) &= 2\pi f_c t + 2\pi k_F \int_0^t A_m \cos(2\pi f_m r) dr \\ &= 2\pi f_c t + 2\pi \frac{k_F A_m}{2\pi f_m} \sin(2\pi f_m t) \\ &= 2\pi f_c t + \frac{k_F A_m}{f_m} \sin(2\pi f_m t) \\ (1) \quad &= 2\pi f_c t + \beta \sin(2\pi f_m t), \quad t \in \mathbb{R} \end{aligned}$$

where

$$\beta := \frac{\Delta f}{f_m} \quad \text{and} \quad \Delta f := k_F A_m.$$

Next,

$$\begin{aligned} \cos(\theta_{\text{FM}}(t)) &= \cos(2\pi f_c t + \beta \sin(2\pi f_m t)) \\ (2) \quad &= \Re(e^{j2\pi f_c t} e^{j\beta \sin(2\pi f_m t)}), \quad t \in \mathbb{R}. \end{aligned}$$

The function $t \rightarrow e^{j\beta \sin(2\pi f_m t)}$ being continuous and periodic with period $T_m = \frac{1}{f_m}$, it admits the Fourier series representation

$$e^{j\beta \sin(2\pi f_m t)} = \sum_k c_k e^{j2\pi k f_m t}, \quad t \in \mathbb{R}$$

with

$$c_k = \frac{1}{T_m} \int_{-\frac{T_m}{2}}^{\frac{T_m}{2}} e^{j\beta \sin(2\pi f_m t)} e^{-j2\pi k f_m t} dt, \quad k = 0, \pm 1, \pm 2, \dots$$

Now fix $k = 0, \pm 1, \pm 2, \dots$. Upon making the change of variable $x = 2\pi f_m t$, we get

$$\begin{aligned} c_k &= \frac{1}{T_m} \int_{-\frac{T_m}{2}}^{\frac{T_m}{2}} e^{j\beta \sin(2\pi f_m t)} e^{-j2\pi k f_m t} dt \\ &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(\beta \sin(x) - kx)} dx \\ (3) \quad &= J_k(\beta) \end{aligned}$$

where

$$J_k(\beta) := \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(\beta \sin(x) - kx)} dx, \quad \beta \in \mathbb{R}$$

defines the k^{th} order Bessel function of the first kind.

Substituting we find

$$e^{j\beta \sin(2\pi f_m t)} = \sum_k J_k(\beta) e^{j2\pi k f_m t}, \quad t \in \mathbb{R}.$$

Therefore,

$$\begin{aligned} A_c \cos(\theta_{\text{FM}}(t)) &= A_c \Re \left(e^{j2\pi f_c t} e^{j\beta \sin(2\pi f_m t)} \right) \\ &= A_c \Re \left(e^{j2\pi f_c t} \sum_k J_k(\beta) e^{j2\pi k f_m t} \right) \\ &= A_c \sum_k J_k(\beta) \Re \left(e^{j2\pi f_c t} e^{j2\pi k f_m t} \right) \\ (4) \quad &= A_c \sum_k J_k(\beta) \cos(2\pi (f_c + k f_m) t), \quad t \in \mathbb{R}. \end{aligned}$$

In the frequency domain this last relationship becomes

$$(5) \quad \begin{aligned} S_{\text{FM}}(f) \\ &= \frac{A_c}{2} \sum_k J_k(\beta) (\delta(f - (f_c + k f_m)) + \delta(f + (f_c + k f_m))) \end{aligned}$$

for all f in \mathbb{R} . Thus, although the single-tone signal m has frequency content *only* at the frequencies $f = \pm f_m$, the corresponding FM wave has *infinite* bandwidth since it displays frequency content at the countably infinite set of frequencies

$$f = \pm(f_c + k f_m), \quad k = 0, \pm 1, \dots$$

Narrow-band vs wide-band FM

Using elementary trigonometric formulae, we observe

$$\begin{aligned} s_{\text{FM}}(t) &= A_c \cos(\theta_{\text{FM}}(t)) \\ &= A_c \cos \left(2\pi f_c t + 2\pi k_F \int_0^t m(r) dr \right) \\ &= A_c \cos(2\pi f_c t) \cos \left(2\pi k_F \int_0^t m(r) dr \right) \\ (6) \quad &\quad - A_c \sin(2\pi f_c t) \sin \left(2\pi k_F \int_0^t m(r) dr \right), \quad t \in \mathbb{R} \end{aligned}$$

Narrow-band FM is characterized by

$$(7) \quad 2\pi k_F \left| \int_0^t m(r) dr \right| \ll 1, \quad t \in \mathbb{R}$$

in which case

$$\sin \left(2\pi k_F \int_0^t m(r) dr \right) \simeq 2\pi k_F \int_0^t m(r) dr$$

and

$$\cos \left(2\pi k_F \int_0^t m(r) dr \right) \simeq 1$$

for all t in \mathbb{R} . Therefore, we have the approximation

$$(8) \quad s_{\text{FM}}(t) \simeq s_{\text{NB-FM}}(t), \quad t \in \mathbb{R}$$

where the narrow-band FM signal $s_{\text{NB-FM}} : \mathbb{R} \rightarrow \mathbb{R}$ is defined by

$$(9) \quad \begin{aligned} s_{\text{NB-FM}}(t) &= A_c \cos(2\pi f_c t) \\ &- A_c \sin(2\pi f_c t) \left(2\pi k_F \int_0^t m(r) dr \right), \quad t \in \mathbb{R}. \end{aligned}$$

In other words, when condition (7) holds, the FM waveform s_{FM} is well approximated by $s_{\text{NB-FM}}$ and therefore can be replaced by it. The advantage of doing so is that the signal $s_{\text{NB-FM}}$ is AM-like in its structure and can be generated easily according to techniques developed for amplitude modulation. *Wide-band FM* arises when the condition (7) fails to hold.

Carson's formula

The realization that the spectrum of s_{FM} has infinite extent leads to the following practical concern: How much bandwidth is needed to transmit s_{FM} without too much distortion?

One answer to this question was given by Carson, and is summarized in the formula that carries his name: Carson's formula states that the transmission bandwidth B_T of the FM wave associated with the single-tone signal m should be set to

$$(10) \quad \begin{aligned} B_{T,\text{Carson}} &:= 2f_m + 2\Delta f \\ &= 2f_m (1 + \beta) \end{aligned}$$

since $\Delta f = f_m \beta$ by definition.

One way to generalize Carson's bandwidth formula could proceed by *formally* giving the quantities f_m and β interpretations which do not rely on the specific form of the information-bearing signal m . We do this as follows:

In the single-tone case, the frequency f_m can be interpreted as the cutoff frequency of the signal – In other words, f_m is the bandwidth of the signal. On the other hand, Δf can be viewed as describing the largest possible excursion of the instantaneous frequency from f_c : Indeed, the instantaneous frequency of the FM wave at time t is given by

$$\frac{1}{2\pi} \frac{d}{dt} \theta_{\text{FM}}(t) = f_c + k_F A_m \cos(2\pi f_m t)$$

and the corresponding deviation in instantaneous frequency at time t is simply

$$\frac{1}{2\pi} \frac{d}{dt} \theta_{\text{FM}}(t) - f_c = k_F A_m \cos(2\pi f_m t).$$

Therefore, the maximal deviation from f_c is given by

$$\sup (|k_F A_m \cos(2\pi f_m t)|, \quad t \in \mathbb{R}) = k_F A_m = \Delta f.$$

Now consider an information bearing signal which is bandlimited with cutoff frequency $W > 0$. With the discussion for the single-tone modulating signal in mind, it is natural to replace in Carson's formula f_m by W and Δf by D with

$$D := \sup (k_F |m(t)|, \quad t \in \mathbb{R}).$$

This suggests the approximation

$$B_T \simeq B_{T, \text{Carson}}$$

with

$$(11) \quad \begin{aligned} B_{T, \text{Carson}} &:= 2W + 2D \\ &= 2W(1 + \beta) \end{aligned}$$

where β is defined as

$$\beta := \frac{D}{W} = \frac{\sup (k_F |m(t)|, \quad t \in \mathbb{R})}{W}.$$

At this point, you may feel that the generalized Carson's formula discussed above is simply a formal expression without much practical grounding. We now show through an approximation argument (see below) that the bandwidth as given by $B_{T,\text{Carson}}$ is indeed meaningful from an engineering point of view.

The basic idea is to characterize the spectrum of the FM wave associated with a *sampled* version of the information-bearing signal. Thus, fix $T > 0$. We approximate the information-bearing signal $m : \mathbb{R} \rightarrow \mathbb{R}$ by the staircase approximation $m_T^* : \mathbb{R} \rightarrow \mathbb{R}$ given by

$$m_T^*(t) = m(kT), \quad kT \leq t < (k+1)T$$

with $k = 0, \pm 1, \dots$. We then replace $\theta_{\text{FM}} : \mathbb{R} \rightarrow \mathbb{R}$ as defined above by $\theta_{\text{FM},T}^* : \mathbb{R} \rightarrow \mathbb{R}$ given by

$$\theta_{\text{FM},T}^*(t) = 2\pi f_c t + 2\pi k_F \int_0^t m_T^*(r) dr, \quad t \in \mathbb{R}$$

and write

$$s_{\text{FM},T}^*(t) = A_c \cos(\theta_{\text{FM},T}^*(t)), \quad t \in \mathbb{R}.$$

Fix f in \mathbb{R} . Note that

$$\begin{aligned} S_{\text{FM},T}^*(f) &= \int_{\mathbb{R}} A_c \cos(\theta_{\text{FM},T}^*(t)) e^{-j2\pi f t} dt \\ (12) \quad &= A_c \sum_k \int_{kT}^{(k+1)T} \cos(\theta_{\text{FM},T}^*(t)) e^{-j2\pi f t} dt. \end{aligned}$$

Now, for $k = 0, 1, \dots$, with $kT \leq t < (k+1)T$, we have

$$\begin{aligned} \theta_{\text{FM},T}^*(t) &= 2\pi f_c t + 2\pi k_F \int_0^t m_T^*(r) dr \\ &= 2\pi f_c t + 2\pi k_F \left(T \sum_{\ell=0}^{k-1} m(\ell T) + m(kT)(t - kT) \right) \\ &= 2\pi(f_c + k_F m(kT))(t - kT) + 2\pi T \left(k f_c + k_F \sum_{\ell=0}^{k-1} m(\ell T) \right) \\ (13) \quad &= 2\pi(f_c + k_F m(kT))(t - kT) + 2\pi \gamma_k T \end{aligned}$$

where we have set

$$\gamma_k := kf_c + k_F \left(\sum_{\ell=0}^{k-1} m(\ell T) \right).$$

Direct substitution yields

$$\begin{aligned} & \int_{kT}^{(k+1)T} \cos(\theta_{\text{FM},T}^*(t)) e^{-j2\pi ft} dt \\ &= \int_{kT}^{(k+1)T} \cos(2\pi(f_c + k_F m(kT))(t - kT) + 2\pi\gamma_k T) e^{-j2\pi ft} dt \\ (14) \quad &= e^{-j2\pi k f T} \cdot \int_0^T \cos(2\pi(f_c + k_F m(kT))\tau + 2\pi\gamma_k T) e^{-j2\pi f\tau} d\tau. \end{aligned}$$

To evaluate this last integral, we note that

$$\begin{aligned} & \int_0^T e^{\pm j2\pi((f_c + k_F m(kT))\tau + \gamma_k T)} e^{-j2\pi f\tau} d\tau \\ &= e^{\pm j2\pi\gamma_k T} \int_0^T e^{j2\pi(\pm(f_c + k_F m(kT)) - f)\tau} d\tau \\ &= e^{\pm j2\pi\gamma_k T} \cdot \frac{e^{j2\pi(\pm(f_c + k_F m(kT)) - f)T} - 1}{j2\pi(\pm(f_c + k_F m(kT)) - f)} \\ &= a_k^\pm(f) \frac{\sin(\pi(\pm(f_c + k_F m(kT)) - f)T)}{\pi(\pm(f_c + k_F m(kT)) - f)} \\ (15) \quad &= a_k^\pm(f) \frac{\sin(\pi(f \mp (f_c + k_F m(kT)))T)}{\pi(f \mp (f_c + k_F m(kT)))} \end{aligned}$$

with

$$a_k^\pm(f) = e^{j2\pi\delta_k^\pm(f)T}$$

where

$$\delta_k^\pm(f) = \pm\gamma_k + \frac{1}{2}(\pm(f_c + k_F m(kT)) - f).$$

Recall that the sinc function $\text{sinc} : \mathbb{R} \rightarrow \mathbb{R}$ is given by

$$\text{sinc}(x) = \frac{\sin(\pi x)}{\pi x}, \quad x \in \mathbb{R}.$$

Therefore, for each $k = 0, 1, \dots$, we have

$$\int_{kT}^{(k+1)T} \cos(\theta_{\text{FM},T}^*(t)) e^{-j2\pi ft} dt$$

$$\begin{aligned}
&= \frac{1}{2} a_k^+(f) \frac{\sin(\pi(f - (f_c + k_F m(kT))) T)}{\pi(f - (f_c + k_F m(kT)))} \\
&\quad + \frac{1}{2} a_k^-(f) \frac{\sin(\pi(f + (f_c + k_F m(kT))) T)}{\pi(f + (f_c + k_F m(kT)))} \\
&= \frac{1}{2} a_k^+(f) \cdot \text{sinc}((f - (f_c + k_F m(kT))) T) \\
(16) \quad &\quad + \frac{1}{2} a_k^-(f) \cdot \text{sinc}((f + (f_c + k_F m(kT))) T),
\end{aligned}$$

and we can conclude

$$\begin{aligned}
&\int_0^\infty \cos(\theta_{\text{FM},T}^*(t)) e^{-j2\pi ft} dt \\
&= \frac{1}{2} \sum_{k=0}^\infty a_k^+(f) \cdot \text{sinc}((f - (f_c + k_F m(kT))) T) \\
(17) \quad &\quad + \frac{1}{2} \sum_{k=0}^\infty a_k^-(f) \cdot \text{sinc}((f + (f_c + k_F m(kT))) T).
\end{aligned}$$

The zeroes of the sinc function occur at $x = \pm\ell$, $\ell = 1, 2, \dots$, and its main lobe occupies the interval $[-1, 1]$. As a result, for each $k = 0, 1, \dots$, the main contribution of the term

$$\frac{1}{2} a_k^\pm(f) \cdot \text{sinc}((f \mp (f_c + k_F m(kT))) T)$$

is taking place on an interval centered at

$$\pm(f_c + k_F m(kT))$$

and of length $2/T$, namely

$$\left[\pm(f_c + k_F m(kT)) - \frac{1}{T}, \pm(f_c + k_F m(kT)) + \frac{1}{T} \right].$$

Similar arguments could be made for the case $k = -1, -2, \dots$ and would lead to a similar expression for

$$\int_{-\infty}^0 \cos(\theta_{\text{FM},T}^*(t)) e^{-j2\pi ft} dt, \quad f \in \mathbb{R}.$$

The discussion suggests that most of the spectral content is contained in the interval

$$\left[\pm(f_c - D) - \frac{1}{T}, \pm(f_c + D) + \frac{1}{T} \right]$$

since

$$|k_F m(kT)| \leq D, \quad k = 0, \pm 1, \dots$$

by the definition of D . This leads to estimating the transmission bandwidth of $s_{\text{FM},T}^*$ as being

$$B_T \simeq 2D + \frac{2}{T}.$$

If we sample at the Nyquist rate, that is $T = \frac{1}{2W}$, then the information contained in m is recoverable from m_T^* , and the transmission bandwidths of their corresponding FM waveforms should be commensurate. In short,

$$B^* = 2D + 4W$$

is expected to provide a reasonably good approximation to B_T . Note that

$$B^* = 2D + 2W + 2W = B_{T,\text{Carson}} + 2W$$

so that this argument provides an approximation to the transmission bandwidth of the FM wave s_{FM} which is more conservative than the one provided by Carson's formula. This can be traced to the fact that the approximation is based on a sampling argument.

Immunity of angle modulation to non-linearities

Consider a non-linear device $\varphi : \mathbb{R} \rightarrow \mathbb{R}$ of the form

$$\varphi(x) = \sum_{m=1}^M a_m x^m, \quad x \in \mathbb{R}$$

for some integer $M \geq 2$ and assume $a_M \neq 0$.

For each t in \mathbb{R} , we note that

$$\varphi(s_{\text{FM}}(t)) = \sum_{m=1}^M a_m (A_c \cos(\theta_{\text{FM}}(t)))^m$$

$$\begin{aligned}
&= \sum_{m=1}^M a_m A_c^m (\cos(\theta_{\text{FM}}(t)))^m \\
(18) \quad &= \sum_{m=1}^M a_m A_c^m \left(\sum_{k=0}^m a_{m,k} \cos(k\theta_{\text{FM}}(t)) \right)
\end{aligned}$$

as we invoke Lemma 0.1 in the last step. Interchanging the order of summation we conclude that

$$\begin{aligned}
\varphi(s_{\text{FM}}(t)) &= \sum_{m=1}^M a_m A_c^m a_{m,0} \\
&\quad + \sum_{k=1}^M \left(\sum_{m=k}^M a_m A_c^m a_{m,k} \right) \cos(k\theta_{\text{FM}}(t)) \\
(19) \quad &= \sum_{\ell=0}^M B_{M,\ell} \cos(\ell\theta_{\text{FM}}(t))
\end{aligned}$$

with

$$(20) \quad B_{M,\ell} = \begin{cases} \sum_{m=1}^M a_m A_c^m a_{m,0} & \text{if } \ell = 0 \\ \sum_{m=\ell}^M a_m A_c^m a_{m,\ell} & \text{if } \ell = 1, \dots, M. \end{cases}$$

For each $\ell = 1, \dots, M$, the signal $t \rightarrow \cos(\ell\theta_{\text{FM}}(t))$ is the FM waveform at carrier frequency ℓf_c generated by the signal $t \rightarrow \ell m(t)$. According to the generalized Carson's rule, for all practical intent, we can view this signal as a bandpass signal whose (transmission) bandwidth B_ℓ is given by

$$B_\ell = 2(W + D_\ell) \quad \text{with} \quad D_\ell = \ell D$$

since

$$\begin{aligned}
D_\ell &= \sup(k_F |\ell m(t)|, t \in \mathbb{R}) \\
(21) \quad &= \ell \sup(k_F |m(t)|, t \in \mathbb{R}) = \ell D.
\end{aligned}$$

Under the appropriate conditions each of the components $t \rightarrow \cos(\ell\theta_{\text{FM}}(t))$ can be extracted from $\varphi(s_{\text{FM}})$ by means of bandpass filtering. For instance, to recover s_{FM} from $\varphi(s_{\text{FM}})$ we pass the latter through a bandpass filter centered at f_c with bandwidth B_1 such that

$$f_c + \frac{B_1}{2} < 2f_c - \frac{B_2}{2}.$$

This is equivalent to

$$f_c + (W + D) < 2f_c - (W + 2D),$$

and requires that the condition

$$2W + 3D < f_c$$

holds.

Similar arguments can be given for extracting $t \rightarrow \cos(\theta_{\text{FM}}(t))$ by means of bandpass filtering.

Generating FM signals

Indirect method of Armstrong We seek to generate the FM signal $s_{\text{FM}} : \mathbb{R} \rightarrow \mathbb{R}$ associated with the information-bearing signal m , say

$$s_{\text{FM}}(t) = A_c \cos(\theta_{\text{FM}}(t)), \quad t \in \mathbb{R}$$

with

$$\theta_{\text{FM}}(t) = 2\pi f_c t + 2\pi k_F \int_0^t m(r) dr, \quad t \in \mathbb{R}$$

for some given $k_F > 0$. We are in the situation when the condition (7) fails to hold for the choice of k_F so that $s_{\text{NB-FM}}$ is not a good approximation to the desired FM signal s_{FM} .

We begin by writing $k_F = Mk_F^*$ for some positive integer M , so that the condition (7) now holds for k_F^* , namely

$$2\pi k_F^* \left| \int_0^t m(r) dr \right| \ll 1, \quad t \in \mathbb{R}.$$

Under this condition the FM signal $s_{\text{FM}}^* : \mathbb{R} \rightarrow \mathbb{R}$ given by

$$s_{\text{FM}}^* = A_c \cos \left(2\pi f_c t + 2\pi k_F^* \int_0^t m(s) ds \right), \quad t \in \mathbb{R}$$

can be well approximated by the narrow-band FM signal $s_{\text{NB-FM}}^* : \mathbb{R} \rightarrow \mathbb{R}$ defined by

$$(22) \quad \begin{aligned} s_{\text{NB-FM}}^*(t) &= A_c \cos(2\pi f_c t) \\ &\quad - A_c \sin(2\pi f_c t) \left(2\pi k_F^* \int_0^t m(r) dr \right), \quad t \in \mathbb{R}. \end{aligned}$$

Next, the narrow-band FM signal $s_{\text{NB-FM}}^* : \mathbb{R} \rightarrow \mathbb{R}$ is converted to the desired wide-band FM signal as follows: Consider a non-linear device $\varphi : \mathbb{R} \rightarrow \mathbb{R}$ of the form

$$\varphi(x) = \sum_{m=1}^M a_m x^m, \quad x \in \mathbb{R}$$

with $a_M \neq 0$.

For each t in \mathbb{R} , with

$$\theta_{\text{FM}}^*(t) = 2\pi f_c t + 2\pi k_F^* \int_0^t m(r) dr,$$

we note from (19)-(20) that

$$(23) \quad \varphi(s_{\text{FM}}^*(t)) = \sum_{\ell=0}^N B_{N,\ell} \cos(\ell \theta_{\text{FM}}^*(t))$$

with the coefficients as given by (20).

By the same arguments as given earlier in the discussion of immunity of angle modulation to non-linearities, we can extract the signal $t \rightarrow \cos(M\theta_{\text{FM}}^*(t))$ by feeding the signal $t \rightarrow \varphi(s_{\text{FM}}^*(t))$ through a bandpass filter with center frequency Nf_c and bandwidth B_M^* given by

$$B_M^* = 2(W + D_M^*)$$

where for each $\ell = 1, 2, \dots$, we have

$$(24) \quad \begin{aligned} D_\ell^* &= \sup (k_F^* |\ell m(t)|, t \in \mathbb{R}) \\ &= \ell \sup (k_F^* |m(t)|, t \in \mathbb{R}) \\ &= \frac{\ell}{M} \sup (k_F |m(t)|, t \in \mathbb{R}). \\ &= \frac{\ell}{M} \cdot D. \end{aligned}$$

As a result,

$$B_M^* = 2(W + D_M^*) = 2(W + D)$$

as should be expected!

Direct method Using a voltage-controlled oscillator (VCO)

Demodulation of FM signals

The FM waveform is given by

$$s_{\text{FM}}(t) = A_c \cos(\theta_{\text{FM}}(t)), \quad t \in \mathbb{R}$$

with

$$\theta_{\text{FM}}(t) = 2\pi f_c t + 2\pi k_F \int_0^t m(r) dr, \quad t \in \mathbb{R}.$$

Assuming sufficient differentiability for m , we note that

$$\begin{aligned} \frac{d}{dt} s_{\text{FM}}(t) &= -A_c \left(\frac{d}{dt} \theta_{\text{FM}}(t) \right) \cdot \sin(\theta_{\text{FM}}(t)) \\ &= -A_c (2\pi f_c + 2\pi k_F m(t)) \cdot \sin(\theta_{\text{FM}}(t)) \\ (25) \quad &= -2\pi A_c (f_c + k_F m(t)) \cdot \sin(\theta_{\text{FM}}(t)), \quad t \in \mathbb{R}. \end{aligned}$$

This calculation highlights the fact that differentiating an FM waveform produces a signal that combines both amplitude and angle modulation. It raises the possibility of using an envelope detector to extract the information bearing signal m . This will be possible if

$$f_c + k_F m(t) > 0, \quad t \in \mathbb{R}.$$

This occurs when

$$D < f_c.$$

The analysis just given is predicated on the amplitude of the FM waveform remaining constant over time. In practice, this condition is not expected to hold. In fact in a number of situations, amplitude distortion can be significant and it is appropriate to model the received signal $s_{\text{FM,Rec}} : \mathbb{R} \rightarrow \mathbb{R}$ to be of the form

$$(26) \quad s_{\text{FM,Rec}}(t) = A(t) \cos(\theta_{\text{FM}}(t)), \quad t \in \mathbb{R}$$

for some $A : \mathbb{R} \rightarrow \mathbb{R}$ with

$$(27) \quad A(t) > 0, \quad t \in \mathbb{R}.$$

Under (26)-(27) the earlier procedure of differentiating the incoming signal and passing the result through an envelope detector will not work anymore: Indeed, assuming enough differentiability, we now have

$$\frac{d}{dt} s_{\text{FM,Rec}}(t)$$

$$\begin{aligned}
&= -A(t) \left(\frac{d}{dt} \theta_{\text{FM}}(t) \right) \cdot \sin(\theta_{\text{FM}}(t)) + \left(\frac{d}{dt} A(t) \right) \cdot \cos(\theta_{\text{FM}}(t)) \\
(28) \quad &= -2\pi A(t) (f_c + k_F m(t)) \cdot \sin(\theta_{\text{FM}}(t)) + \left(\frac{d}{dt} A(t) \right) \cdot \cos(\theta_{\text{FM}}(t)).
\end{aligned}$$

The approach based on envelope detection used when the amplitude remained constant will not work here due to the presence of the *unknown* and *time-varying* term

$$\left(\frac{d}{dt} A(t) \right) \cdot \cos(\theta_{\text{FM}}(t)), \quad t \in \mathbb{R}.$$

We can remedy to this difficulty by preprocessing $s_{\text{FM,Rec}}$ with the aim of extracting the original waveform s_{FM} . One possible way to achieve this goal is presented next.

Consider the hard-limiter $\Phi : \mathbb{R} \rightarrow \mathbb{R}$ given by

$$(29) \quad \Phi(x) = \begin{cases} -1 & \text{if } x < 0 \\ 0 & \text{if } x = 0 \\ 1 & \text{if } x > 0 \end{cases}$$

For each t in \mathbb{R} , as we recall that $A(t) > 0$, we note that

$$\begin{aligned}
(30) \quad v(t) &:= \Phi(s_{\text{FM-Rec}}(t)) \\
&= \Phi(A(t) \cos(\theta_{\text{FM}}(t))) \\
&= \Phi(\cos(\theta_{\text{FM}}(t))).
\end{aligned}$$

Next, observe that the mapping $\theta \rightarrow \Phi(\cos \theta)$ is a periodic function with period 2π – In fact, this function is just the periodic square wave function and therefore admits a Fourier series representation, say

$$(31) \quad \Phi(\cos \theta) = \sum_k c_k e^{jk\theta}, \quad \theta \in \mathbb{R}$$

with

$$c_k = \frac{1}{2\pi} \int_{-\pi}^{\pi} \Phi(\cos \theta) e^{-jk\theta} d\theta, \quad k = 0, \pm 1, \dots$$

After some straightforward calculations (see below) we conclude that

$$\Phi(\cos \theta) = \frac{4}{\pi} \left(\sum_{\ell=0}^{\infty} \frac{(-1)^\ell}{2\ell+1} \cos((2\ell+1)\theta) \right), \quad \theta \in \mathbb{R}.$$

As a result,

$$\begin{aligned}
 v(t) &= \Phi(\cos(\theta_{\text{FM}}(t))) \\
 (32) \quad &= \frac{4}{\pi} \left(\sum_{\ell=0}^{\infty} \frac{(-1)^\ell}{2\ell+1} \cos((2\ell+1)\theta_{\text{FM}}(t)) \right).
 \end{aligned}$$

Again, as was the case in the discussion of demodulation of FM signals, we note that for each $\ell = 0, 1, \dots$, the signal $t \rightarrow \cos((2\ell+1)\theta_{\text{FM}}(t))$ is the FM waveform at carrier frequency $(2\ell+1)f_c$ generated by the signal $t \rightarrow (2\ell+1)m(t)$. According to the generalized Carson's rule, we can view this signal as a bandpass signal with bandwidth $B_{2\ell+1}$. If we pass the signal v through a bandpass filter with center frequency f_c and bandwidth $2W + 2D$ we will collect the signal $t \rightarrow \frac{4}{\pi} \cos(\theta_{\text{FM}}(t))$, as required. The earlier procedure outlined earlier, namely feeding into a differentiator followed by an envelope detector, can now be used on this resulting waveform.

Properties of Bessel functions

0. For each $k = 0, \pm 1, \dots$ and every β in \mathbb{R} , $J_k(\beta)$ is an element of \mathbb{R} .

Proof. Fix $k = 0, \pm 1, \dots$ and β in \mathbb{R} . Note that

$$\begin{aligned}
 J_k(\beta)^* &= \left(\frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(\beta \sin x - kx)} dx \right)^* \\
 &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-j(\beta \sin x - kx)} dx \\
 &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(-\beta \sin x + kx)} dx \\
 &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(\beta \sin(-x) - k(-x))} dx \\
 &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(\beta \sin y - ky)} dy \\
 (33) \qquad &= J_k(\beta).
 \end{aligned}$$

Hence $J_k(\beta)^* = J_k(\beta)$, and $J_k(\beta)$ is an element of \mathbb{R} . ■

1. For each $k = 0, 1, \dots$, we have

$$J_{-k}(\beta) = (-1)^k J_k(\beta), \quad \beta \in \mathbb{R}.$$

Proof. Fix $k = 0, 1, \dots$ and β in \mathbb{R} . Using the change of variable $y = \pi - x$ we find

$$\begin{aligned}
 J_{-k}(\beta) &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(\beta \sin x + kx)} dx \\
 &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(\beta \sin(\pi - y) + k(\pi - y))} dy \\
 &= \left(\frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(\beta \sin y - y)} dy \right) \cdot e^{jk\pi} \\
 (34) \qquad &= (-1)^k J_k(\beta)
 \end{aligned}$$

since $e^{jk\pi} = (-1)^k$. ■

2. For each $k = 0, 1, \dots$, we have

$$J_k(-\beta) = (-1)^k J_k(\beta), \quad \beta \geq 0.$$

Proof. Fix $k = 0, 1, \dots$ and $\beta \geq 0$. We note that

$$\begin{aligned} J_k(-\beta) &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(-\beta \sin y - ky)} dy \\ &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(\beta \sin(-y) + k(-y))} dy \\ &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j(\beta \sin x + kx)} dx \\ (35) \qquad &= J_{-k}(\beta) \end{aligned}$$

and the conclusion follows by Fact 1. ■

3. We have

$$J_0(\beta) = 1 + O(\beta) \quad (\beta \rightarrow 0).$$

Proof. Fix β in \mathbb{R} . From the definitions we see that

$$\begin{aligned} J_0(\beta) - 1 &= \frac{1}{2\pi} \int_{-\pi}^{\pi} (e^{j\beta \sin x} - 1) dx \\ (36) \qquad &= \frac{1}{2\pi} \int_{-\pi}^{\pi} \left(\int_0^{\beta \sin x} j e^{jt} dt \right) dx \end{aligned}$$

so that

$$\begin{aligned} |J_0(\beta) - 1| &\leq \frac{1}{2\pi} \int_{-\pi}^{\pi} \left| \int_0^{\beta \sin x} j e^{jt} dt \right| dx \\ &\leq \frac{1}{2\pi} \int_{-\pi}^{\pi} \left| \int_0^{|\beta \sin x|} |j e^{jt}| dt \right| dx \\ &\leq \frac{1}{2\pi} \int_{-\pi}^{\pi} |\beta| |\sin x| dx \\ (37) \qquad &\leq |\beta| \end{aligned}$$

and the conclusion follows. ■

4. We have

$$J_1(\beta) = \frac{\beta}{2}(1 + o(1)) \quad (\beta \rightarrow 0).$$

5. For each $\ell = 0, 1, \dots$ we have

$$J_\ell(\beta) = \frac{\beta^\ell}{2^\ell \ell!}(1 + o(1)) \quad (\beta \rightarrow 0).$$

6. For each β in \mathbb{R} , we have

$$\sum_{\ell} |J_\ell(\beta)|^2 = 1.$$

Proof. For each β in \mathbb{R} , the function $x \rightarrow e^{j\beta \sin x}$ is periodic with period 2π and therefore admits a Fourier series representation. It is a simple matter to see that

$$e^{j\beta \sin x} = \sum_{\ell} J_\ell(\beta) e^{j\ell x}$$

and by Parseval's Theorem we get

$$\frac{1}{2\pi} \int_{-\pi}^{\pi} |e^{j\beta \sin x}|^2 dx = \sum_{\ell} |J_\ell(\beta)|^2.$$

The conclusion follows from the fact that

$$|e^{j\beta \sin x}|^2 = 1, \quad x \in \mathbb{R}.$$

■

On powers of $\cos \theta$

Given is θ in \mathbb{R} . We are interested in understanding how to compute

$$(\cos \theta)^m, \quad m = 1, 2, \dots$$

We shall repeatedly use the trigonometric identity

$$2 \cos \alpha \cos \beta = \cos(\alpha + \beta) + \cos(\alpha - \beta)$$

for arbitrary α and β in \mathbb{R} .

For $m = 2$, we have

$$(38) \quad (\cos \theta)^2 = \frac{\cos(2\theta) + 1}{2}.$$

Next, with $m = 3$,

$$\begin{aligned} (\cos \theta)^3 &= \frac{\cos(2\theta) + 1}{2} \cdot \cos \theta \\ &= \frac{\cos(2\theta) \cos \theta + \cos \theta}{2} \\ &= \frac{\frac{\cos(3\theta) + \cos \theta}{2} + \cos \theta}{2} \\ (39) \quad &= \frac{\cos(3\theta) + 3 \cos \theta}{4} \end{aligned}$$

Building on the pattern emerging from these calculations we now set out to prove the following fact.

Lemma 0.1 *Given θ in \mathbb{R} , for each $m = 1, 2, \dots$, there exist scalars $a_{m,0}, \dots, a_{m,m}$, independent of θ , such that*

$$(40) \quad (\cos \theta)^m = \sum_{k=0}^m a_{m,k} \cos(k\theta).$$

Proof. The proof proceeds by induction. The conclusion (40) is true for $m = 1$ (with $a_{1,0} = 0$ and $a_{1,1} = 1$), for $m = 2$ (with $a_{2,0} = \frac{1}{2}$, $a_{2,1} = 0$ and $a_{2,2} = \frac{1}{2}$) and for $m = 3$ (with $a_{3,0} = 0$, $a_{3,1} = \frac{3}{2}$, $a_{3,2} = 0$ and $a_{3,3} = \frac{1}{4}$).

Now assume (40) to hold for some $m \geq 2$. We note that

$$\begin{aligned}
 (\cos \theta)^{m+1} &= (\cos \theta)^m \cdot \cos \theta \\
 &= \left(\sum_{k=0}^m a_{m,k} \cos(k\theta) \right) \cdot \cos \theta \\
 &= a_{m,0} \cos \theta + \sum_{k=1}^m a_{m,k} \cos(k\theta) \cos \theta \\
 &= a_{m,0} \cos \theta + \sum_{k=1}^m a_{m,k} \frac{\cos((k+1)\theta) + \cos((k-1)\theta)}{2} \\
 &= a_{m,0} \cos \theta + \frac{1}{2} \sum_{k=1}^m a_{m,k} \cos((k+1)\theta) + \frac{1}{2} \sum_{k=1}^m a_{m,k} \cos((k-1)\theta) \\
 &= a_{m,0} \cos \theta + \frac{1}{2} \sum_{k=2}^{m+1} a_{m,k-1} \cos(k\theta) + \frac{1}{2} \sum_{k=0}^{m-1} a_{m,k+1} \cos(k\theta) \\
 (41) \quad &= \sum_{k=0}^{m+1} a_{m+1,k} \cos(k\theta)
 \end{aligned}$$

with

$$a_{m+1,k} = \begin{cases} \frac{a_{m,1}}{2} & \text{if } k = 0 \\ a_{m,0} + \frac{a_{m,2}}{2} & \text{if } k = 1 \\ \frac{1}{2} (a_{m,k-1} + a_{m,k+1}) & \text{if } k = 2, \dots, m-1 \\ \frac{a_{m,m-1}}{2} & \text{if } k = m \\ \frac{a_{m,m}}{2} & \text{if } k = m+1 \end{cases}$$

by direct inspection. This completes the proof of Lemma 0.1. ■

Computing the Fourier coefficients for $\Phi(\cos \theta)$

Recall that

$$c_k = \frac{1}{2\pi} \int_{-\pi}^{\pi} \Phi(\cos \theta) e^{-jk\theta} d\theta, \quad k = 0, \pm 1, \dots$$

with Φ as defined at (29). Thus,

$$\begin{aligned}
 & \int_{-\pi}^{\pi} \Phi(\cos \theta) e^{-jk\theta} d\theta, \\
 &= - \int_{-\pi}^{-\frac{\pi}{2}} e^{-jk\theta} d\theta + \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} e^{-jk\theta} d\theta - \int_{\frac{\pi}{2}}^{\pi} e^{-jk\theta} d\theta \\
 &= \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} e^{-jk\theta} d\theta - \int_{\frac{\pi}{2}}^{\pi} (e^{jk\theta} + e^{-jk\theta}) d\theta \\
 &= \frac{e^{-jk\frac{\pi}{2}} - e^{jk\frac{\pi}{2}}}{-jk} - 2 \int_{\frac{\pi}{2}}^{\pi} \cos(k\theta) d\theta \\
 &= \frac{2}{k} \sin\left(k\frac{\pi}{2}\right) - \frac{2}{k} \left(\sin(k\pi) - \sin\left(k\frac{\pi}{2}\right) \right) \\
 &= \frac{4}{k} \sin\left(k\frac{\pi}{2}\right) - \frac{2}{k} \sin(k\pi) \\
 (42) \quad &= \frac{4}{k} \sin\left(k\frac{\pi}{2}\right).
 \end{aligned}$$

It is plain that

$$\sin\left(k\frac{\pi}{2}\right) = \begin{cases} (-1)^\ell & \text{if } k = 2\ell + 1 \\ 0 & \text{if } k = 2\ell \end{cases}$$

with $\ell = 0, \pm 1, \dots$